

# Monte Carlo Methods Based on Analytic Extension of the Resolvent of the Helmholtz Equation<sup>1</sup>

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Let us consider the Dirichlet problem for the Helmholtz equation in a domain  $D \subset \mathbb{R}^3$  with boundary  $\Gamma$ :

$$(\Delta + c)u = -g, \quad u|_{\Gamma} = \varphi \quad (1)$$

Let us assume that the following conditions hold. The function  $g$  is satisfied Holder condition in  $\bar{D}$ ,  $D$  is a bounded open set in  $\mathbb{R}^3$  with a regular boundary  $\Gamma$ , the function  $\varphi$  is continuous on  $\Gamma$ ,  $c < c^*$ , where  $-c^*$  is the first eigen value of Laplace operator defined on the domain  $D$ .

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It is well known that solution to the problem (1) satisfies the integral equation

$$u = \int_D k(r, r'; c) u(r') dr' + h(r), \text{ or } u = K_c u + h, \quad (2)$$

One of approaches to constructing statistical algorithms for solving (2) is based on the following presentation

$$u = (I - K_c)^{-1} h = h + R_c h, \quad R_c = \sum_{i=1}^{\infty} K_c^i \quad (3)$$

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Let us consider the following problem

$$Ku = \lambda u$$

We can use well known formula to calculate first eigenvalue

$$\lambda_1 = \lim_{i \rightarrow \infty} \frac{a_i}{a_{i+1}}$$

$$R_ch = \sum_{i=0}^{\infty} a_i c^i, \quad a_i = K^{i+1} h$$

Using the same methods by parametric derivatives. We consider the parametric derivatives  $u^{(m)}$  of the solution to the particular problem [Mikhailov G.A.]

$$(\Delta + c)u = 0, \quad u|_{\Gamma} = 1$$

It obviously, that

$$(\Delta + c)u^{(1)} = -u, \quad u^{(1)}|_{\Gamma} = 0$$

$$(\Delta + c)u^{(m)} = -mu^{(m-1)}, \quad u^{(m)}|_{\Gamma} = 0, \quad m = 1, 2, \dots$$

Therefore

$$c^* - c = \lim_{m \rightarrow \infty} \frac{mu^{(m-1)}}{u^{(m)}}$$

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It is known (see [Elepov B.S., Mikhailov G.A.(1969)]), that under above stated conditions the probability representation for the solution to problem (1) has the form

$$u(x) = E\left[\int_0^{\tau} e^{s(t;c)} g(\xi_t) dt + e^{s(\tau;c)} \varphi(\xi_{\tau})\right], \quad s(t;c) = \int_0^t c(\xi_{t'}) dt', \quad (4)$$

where  $\xi_t$  is the diffusion process corresponding to the Laplace operator and starting at the point  $x$ ,  $\tau$  is the moment of the first exit of the process from the domain  $D$ .

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Let us consider the remainder of Neumann series

$$Rem_n = \sum_{i=n+1}^{\infty} a_i c^i$$

When  $u(r,c)$  is meromorphic function we have the following estimate [Kublanovskaya V.N.,1959]

$$|Rem_n| = O(\delta^{n+1} n^{r-1})$$

Here  $\delta = |c/c_1^{(k)}|$ ,  $c_1^{(k)}$  is any of the poles lying on the convergence circle,  $r$  is the highest multiplicity of this poles

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# The relative error of the eigenvalue algorithm

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Now note that

$$\varepsilon_i = \left| \left( \frac{a_i}{a_{i+1}} - \lambda_1 \right) / \lambda_1 \right|$$

In this case

$$\varepsilon_i = O(\delta_1^{n+1}),$$

where  $\delta_1 = |\lambda_1/\lambda_2|$ , when  $\lambda_1, \lambda_2$  are the simple poles.

$$\varepsilon_i = O(\delta_1^{n+1} i^{r-1}),$$

when  $\lambda_1$  is the simple pole,  $\lambda_2$  is the multiple pole ( $r$  is the multiplicity of pole).

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$$\varphi : \eta \rightarrow \mathbf{c}, \quad \varphi(\eta) = \sum_{i=1}^{\infty} \mathbf{s}_i \eta^i.$$

$$F(\eta) = u(\varphi(\eta), r) = \sum_{i=0}^{\infty} b_i \eta^i,$$

where

$$b_i = \sum_{k=0}^i a_k d_k^{(i)}, \quad d_k^{(i)} = \frac{1}{i!} \left( \frac{\partial^i}{\partial \eta^i} (\varphi(\eta))^k \right)_{\eta=0}$$

$$u(\mathbf{c}, r) = u(\varphi(\varphi^{-1}(\mathbf{c})), r) = F(\eta) = \sum_{i=0}^{\infty} b_i \eta^i =$$

$$= \sum_{i=0}^{\infty} \eta^i \sum_{k=1}^i d_k^{(i)} a_i = \sum_{k=1}^{\infty} a_i l_i,$$

$$\text{where } l_i = \sum_{k=i}^{\infty} d_k^{(i)} \eta^i$$

$$u(c, r) = \mathbb{E}\left[\int_0^\tau e^{\varphi(\eta)t} g(\xi_t) dt + e^{\varphi(\eta)\tau} \phi(\xi_\tau)\right]$$

$$e^{\varphi(\eta)t} = \sum_{i=0}^{\infty} \frac{\varphi^i(\eta)t^i}{i!} = \sum_{i=0}^{\infty} b_i \eta^i,$$

$$\text{where } b_i = \sum_{k=1}^i d_k^{(i)} \frac{t^k}{k!}$$

$$u(c, r) \approx \sum_{k=0}^{\infty} \mathbb{E}\left[\int_0^\tau \frac{t^k}{k!} \left(\sum_{n=k}^m d_k^{(n)} \eta^n\right) g(\xi_t) dt + e^{\varphi(\eta)\tau} \phi(\xi_\tau)\right]$$



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- Kublanovskaya V.N., Sabelfeld K.K.

$$\mathbf{c} = \varphi(\eta) = \frac{4\alpha\eta}{(1-\eta)^2}$$

- Mikhailov G.A.

$$\mathbf{c} = \varphi(\eta) = \frac{\alpha}{(1-\beta\eta)}$$

- 

$$\mathbf{c} = \varphi(\eta) = \alpha + \beta \arctan \eta$$

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Let us consider the following problem

$$\begin{cases} (\Delta + c)^{p+1} u = -g, \\ (\Delta + c)^k u|_{\Gamma} = \varphi_k, \quad k = 0, \dots, p. \end{cases} \quad (5)$$

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Theorem (Mikhailov G.A., Lukinov V.L., 2004)

Under above stated conditions, the  $p$ -th parametric derivative of the solution to problem (1) with the functional parameters

$$\varphi = \sum_{k=0}^p \frac{(-1)^k (c - c_1)^{p-k}}{p!} \varphi_k, \quad g_1 = \frac{(-1)^p}{p!} g \quad (6)$$

is a solution to problem (1) (all derivatives are calculated at the point  $c_1$ ). Here  $p$ -th parametric derivative satisfies the following problem

$$\begin{cases} (\Delta + c)^{p+1} u = -g \frac{c_0^{j+1}}{j!}, \\ (\Delta + c)^k u|_{\Gamma} = \varphi \frac{c_0^{j+1}}{j!}, \quad k = i-1, \\ (\Delta + c)^k u|_{\Gamma} = 0, \quad k = i-2, \dots, 0. \end{cases} \quad (7)$$

## Theorem (2)

If  $\mathbf{c} < \mathbf{c}^*$  and the first spatial derivatives of the functions  $\{u_k^{(i)}\}$ ,  $i = 1, \dots, p + 1$ , are uniformly bounded in  $\bar{D}$ , then

$$|u(r) - E\tilde{\eta}_{1,\varepsilon}^{(p)}| \leq C_p \varepsilon, \quad r \in D, \quad \varepsilon > 0,$$

where

$$\begin{aligned} \tilde{\eta}_{1,\varepsilon}^{(p)} = & \sum_{i=0}^N \left\{ \left[ \prod_{j=0}^{i-1} s(\mathbf{c}, d_j) \right] g(\rho_i) \frac{d_i^2 G(\rho; \mathbf{c}, d_i)}{2nG(\rho; 0, d_i)} \right\}^{(p)} \\ & + \left\{ \left[ \prod_{j=0}^{N-1} s(\mathbf{c}, d_j) \right] \varphi(r_N, \mathbf{c}) \right\}^{(p)}, \end{aligned} \quad (8)$$

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Here,  $d_j = d(r_j)$ ,  $D(r_j)$  is a ball of radius  $d_j$  with the center at the point  $r_j$ ,  $G(\rho; \mathbf{c}, \mathbf{d})$  is the spherical Green's function; and

$$s(\mathbf{c}, \mathbf{d}) = \begin{cases} d\sqrt{c}/\sin(d\sqrt{c}), & c \geq 0; \\ d\sqrt{|c|}/\sinh(d\sqrt{|c|}), & c \leq 0. \end{cases}$$

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$$\begin{aligned} \tilde{\eta}_{1,\varepsilon}^{(p)} &= \sum_{i=0}^N \left\{ \left[ \prod_{j=0}^{i-1} \frac{\sqrt{c}d_j}{\sin(\sqrt{c}d_j)} \right] g(\nu_i, \omega_i) \frac{d_i^3 \sqrt{c} \sin(\sqrt{c}(d_i - \nu_i))}{6(d_i - \nu_i) \sin(\sqrt{c}d_i)} \right. \\ &\quad \left. + \left\{ \left[ \prod_{j=0}^{N-1} \frac{\sqrt{c}d_j}{\sin(\sqrt{c}d_j)} \right] \varphi(r_N, c) \right\}^{(p)} \right\} \end{aligned}$$

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$$\eta_k = \sum_{p=0}^k \left[ \sum_{i=0}^N \left\{ \left[ \prod_{j=0}^{i-1} \frac{\sqrt{c}d_j}{\sin(\sqrt{c}d_j)} \right] \frac{c_0^{i+1}}{i!} g(\nu_i, \omega_i) \frac{d_i^3 \sqrt{c} \sin(\sqrt{c}(a - d_i))}{6(d_i - \nu_i) \sin(\sqrt{c}d_i)} \right. \right. \\ \left. \left. + \left\{ \left[ \prod_{j=0}^{N-1} \frac{\sqrt{c}d_j}{\sin(\sqrt{c}d_j)} \right] \frac{c_0^{i+1}}{i!} \varphi(r_N, c) \right\}^{(p)} \right] \right]$$

The random variables  $\nu_i$  distributed in the interval  $(0, d_i)$  with the probability density  $6x(1 - x/d_i)d_i^{-2}$  and the unit isotropic vectors  $\omega_i$  are modeled by means of the well-known formulas

[S. M. Ermakov and G. A. Mikhailov, 1982].



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## Lemma

Under above stated conditions, and  $\mathbf{c} < \mathbf{c}^*$ . Then

$$\frac{\partial^p u}{\partial \mathbf{c}^p} = u^{(p)} = \mathbb{E} \left[ \int_0^\tau t^p e^{ct} g(\xi_t) dt + \sum_{l=0}^p C_p^l T^{p-l} e^{cT} \frac{\partial^l}{\partial \mathbf{c}^l} \varphi(\xi_T, \mathbf{c}) \right].$$

On the other hand we have

$$u(c, r) = \mathbb{E}\left[\int_0^\tau e^{ct} g(\xi_t) dt + e^{c\tau} \varphi(\xi_\tau)\right] =$$

$$= \sum_{i=0}^{\infty} \mathbb{E}\left[\int_0^\tau \frac{c^i t^i}{i!} g(\xi_t) dt + \frac{c^i \tau^i}{i!} \varphi(\xi_\tau)\right] = \sum_{i=0}^{\infty} c^i a_i,$$

where  $a_i = \mathbb{E}\left[\int_0^\tau \frac{t^i}{i!} g(\xi_t) dt + \frac{\tau^i}{i!} \varphi(\xi_\tau)\right];$

$$R_n = \sum_{i=n}^{\infty} c^i a_i = O(|c/c^*|^{n+1} n^{r-1}).$$

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The oscillations of the plate in a bounded region  $D \subset R^2$  under action of the random field stress  $\sigma(r) = -g(r)$  describes equation of the form

$$\Delta^2 u + cu = -g, \quad u|_{\Gamma} = \varphi_0, \quad \Delta u|_{\Gamma} = \varphi_1, \quad (9)$$

thus it is possible consider the randomness of the boundary functions  $\varphi_0(r)$  and  $\varphi_1(r)$  [Bolotin V.V.,1979] also.

The solution to this problem is a random field. Required to determine its covariance function  $v(r, r') = E[u(r)u(r')]$ .

It is assumed that  $Eg(r) \equiv E\varphi_0(r) \equiv E\varphi_1(r) \equiv 0$  and, hence,  $Eu(r) \equiv 0$ .

Solution  $u(r, \mathbf{c}, \mathbf{g}, \varphi_0, \varphi_1)$  can be approximately replaced the corresponding partial sum of Loran series. To construct estimates of the unknown parameter derivatives of  $u^{(1)}(r, \mathbf{0}, \mathbf{g}, \varphi_0, \varphi_1)$ ,  $u^{(2)}(r, \mathbf{0}, \mathbf{g}, \varphi_0, \varphi_1)$  we use the following way. Differentiating with respect to the parameter  $\mathbf{c}$  equation (??), we get for  $\mathbf{c} = \mathbf{0}$ :

$$\Delta^2 u^{(1)} = -u, \quad u^{(1)} \Big|_{\Gamma} = \Delta u^{(1)} \Big|_{\Gamma} = 0. \quad (10)$$

Hence it is clear that the parametric derivative  $u^{(1)}(r, \mathbf{0}, \mathbf{g}, \varphi_0, \varphi_1)$  is the solution of the following problem:

$$\begin{cases} \Delta^4 u = -\mathbf{g}, & u|_{\Gamma} = 0, & \Delta u|_{\Gamma} = 0, \\ \Delta^2 u|_{\Gamma} = -\varphi_0, & \Delta^3 u|_{\Gamma} = -\varphi_1. \end{cases} \quad (11)$$

The parametric derivative  $u^{(2)}(r, \mathbf{0}, \mathbf{g}, \varphi_0, \varphi_1)/2$  is the solution of the following problem:

$$\begin{cases} \Delta^6 u = \mathbf{g}, & \Delta^k u|_{\Gamma} = 0, & k = 0, \dots, 3, \\ \Delta^4 u|_{\Gamma} = \varphi_0, & \Delta^5 u|_{\Gamma} = \varphi_1 \end{cases} \quad (12)$$

Estimates for solutions of (??), (??) respectively are as follows:

$$\begin{aligned} \tilde{\eta}_{1,\varepsilon}^{(3)} = & \sum_{i=0}^N \left\{ \sum_{k=0}^3 C_3^k S_i^{(3-k)}(0) \frac{G^{(k)}(\rho; 0, d_i)}{G(\rho; 0, d_i)} \right\} \frac{[-d_i^2 g(\rho_i)]}{24} \\ & + S_N^{(3)}(0) \frac{\varphi_1(r_N)}{6} - S_N^{(2)}(0) \frac{\varphi_0(r_N)}{2}, \end{aligned} \quad (13)$$

$$\begin{aligned} \tilde{\eta}_{1,\varepsilon}^{(5)} = & \left\{ \sum_{i=0}^N \sum_{k=0}^5 C_5^k S_i^{(5-k)}(0) \frac{G^{(k)}(\rho; 0, d_i)}{G(\rho; 0, d_i)} \right\} \frac{d_i^2 g(\rho_i)}{480} \\ & - S_N^{(5)}(0) \frac{\varphi_1(r_N)}{120} + S_N^{(4)}(0) \frac{\varphi_0(r_N)}{24}, \end{aligned} \quad (14)$$

where  $S_i(c) = \prod_{j=0}^{i-1} s(c, d_j)$ .



Representation of estimates shows that it is necessary to find values derivatives of  $S_i^{(k)}(0)$  for  $k = \overline{0, 3}$ . Differentiating with respect to parameter  $\mathbf{c}$  the logarithmic derivative of  $S^{(1)} = S(\ln S)^{(1)}$ , we can easily obtain the following relations

$$S^{(2)} = S^{(1)}(\ln S)^{(1)} + S(\ln S)^{(2)},$$

$$S^{(3)} = S^{(2)}\ln S^{(1)} + 2S^{(1)}(\ln S)^{(2)} + S(\ln S)^{(3)},$$

$$S^{(4)} = S^{(3)}\ln S^{(1)} + 3S^{(2)}(\ln S)^{(2)} + 3S^{(1)}(\ln S)^{(3)} + S(\ln S)^{(4)},$$

$$S^{(5)} = S^{(4)}\ln S^{(1)} + 4S^{(3)}(\ln S)^{(2)} + 6S^{(2)}(\ln S)^{(3)} \\ + 4S^{(1)}(\ln S)^{(4)} + S(\ln S)^{(5)}.$$

Derivatives of  $(\ln S)^{(k)}$ ,  $k = \overline{1, 5}$  at  $c = 0$  is easily obtained using equations

$$s(c, d) = 1 + \frac{d^2}{4}c + \frac{3d^4}{2^6}c^2 + \frac{19d^6}{2^8 3^2}c^3 + \frac{211}{2^{14} 3^2}c^4 + \frac{1217}{2^{16} \cdot 3}c^5$$

$$(\ln S_j)^{(1)} = \sum_{j=0}^{i-1} \frac{s^{(1)}(0, d_j)}{s(0, d_j)} = \sum_{j=0}^{i-1} \frac{d_j^2}{4},$$

$$(\ln S_j)^{(2)} = \sum_{j=0}^{i-1} \left\{ \frac{s^{(2)}(0, d_j)}{s(0, d_j)} - \left( \frac{s^{(1)}(0, d_j)}{s(0, d_j)} \right)^2 \right\} = \sum_{j=0}^{i-1} \frac{d_j^4}{2^5},$$

$$(\ln S_j)^{(3)} = \sum_{j=0}^{i-1} \left\{ \frac{s^{(3)}(0, d_j)s(0, d_j) - s^{(2)}(0, d_j)s^{(1)}(0, d_j)}{s^2(0, d_j)} - \frac{2s^{(1)}(0, d_j)[s^{(2)}(0, d_j)s(0, d_j) - (s^{(1)}(0, d_j))^2]}{s^3(0, d_j)} \right\} =$$

$$(\ln S_j)^{(4)} = \sum_{j=0}^{i-1} \left\{ \frac{s^{(4)}(0, d_j)}{s(0, d_j)} - \frac{4s^{(3)}(0, d_j)s^{(1)}(0, d_j)}{s^2(0, d_j)} \right\}$$

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Using the following asymptotic relations

$$\begin{aligned} \frac{J_0(z\sqrt{c})}{J_0(d\sqrt{c})} &\sim 1 + \frac{d^2 - z^2}{4}c + \frac{3d^4 - 4z^2d^2 + z^4}{2^6}c^2 \\ &+ \frac{19d^6 - 3^3z^2d^4 + 3^2z^4d^2 - z^6}{2^83^2}c^3 \\ &+ \frac{211d^8 - 19 \cdot 2^4d^6z^2 + 3^32^2d^4z^4 - 2^4d^2z^6 + z^8}{2^{14}3^2}c^4 \\ &+ \left[ \frac{3 \cdot 1217d^{10} - 211 \cdot 5^2d^8z^2 + 19 \cdot 2^2 \cdot 5^2d^6z^4}{2^{16} \cdot 3^2 \cdot 5^2} \right. \\ &\left. + \frac{-2^2 \cdot 3 \cdot 5^2d^4z^6 + 5^2d^2z^8 - z^{10}}{2^{16} \cdot 3^2 \cdot 5^2} \right] c^5 \end{aligned}$$

$$N_0^{(2)}(z\sqrt{c}) \sim \frac{2}{\pi} \left[ -\frac{1}{2c^2} - \frac{z^2}{2^3c} - \frac{3z^4}{2^7} + \left( \gamma + \ln \frac{z\sqrt{c}}{2} \right) \frac{z^4}{2^5} \right],$$

$$N_0^{(3)}(z\sqrt{c}) \sim \frac{2}{\pi} \left[ \frac{1}{c^3} + \frac{z^2}{2^3c^2} + \frac{z^4}{2^6c} + \frac{z^6}{2^83^2} - \left( \gamma + \ln \frac{z\sqrt{c}}{2} \right) \frac{z^6}{2^73} \right],$$

$$N_0^{(4)}(z\sqrt{c}) \sim \frac{2}{\pi} \left[ -\frac{3}{c^4} - \frac{z^2}{4c^3} - \frac{z^4}{2^6c^2} - \frac{z^6}{2^83c} - \frac{25z^8}{2^{14}3^2} + \left( \gamma + \ln \frac{z\sqrt{c}}{2} \right) \frac{z^8}{2^{13}3} \right],$$

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we derive that

$$G^{(2)}(\rho; 0, d) = \frac{1}{2^8 \pi} (6z^4 - 16z^2 d^2 + 10d^4 + 4z^4 \ln \frac{d}{z}),$$

$$G^{(3)}(\rho; 0, d) = \frac{1}{2^{10} 3^2 \pi} (175d^6 - 270d^4 z^2 + 108d^2 z^4 - 13z^6 - \dots)$$

$$G^{(4)}(\rho; 0, d) = \frac{1}{2\pi} \left( \frac{677d^8}{2^{13}3} + \frac{25z^8}{2^{13}3^2} - \frac{z^6 d^2}{2^7 3} - \frac{23d^6 z^2}{2^6 3^2} + \frac{15z^4 d^4}{2^{10}} \right)$$

$$G^{(5)}(\rho; 0, d) = \frac{1}{2\pi} \left( \frac{71 \cdot 103 d^{10}}{2^{12} \cdot 3 \cdot 5^2} - \frac{137 z^{10}}{2^{15} \cdot 3^2 \cdot 5^2} + \frac{5 z^8 d^2}{2^{13} 3} - \frac{5 \cdot 67}{2^{10}} \right. \\ \left. - \frac{5^2 z^6 d^4}{2^{12} 3} + \frac{5 \cdot 23 z^4 d^6}{2^{10} 3^2} - \frac{z^{10}}{2^{13} \cdot 3 \cdot 5} \ln \frac{d}{z} \right).$$

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$$\begin{aligned} \text{Cov}(u(r_1, c, g, \varphi_0, \varphi_1), u(r_2, c, g, \varphi_0, \varphi_1)) &\approx v(r_1, r_2) + \\ &+ E[cu(r_1, 0, g, \varphi_0, \varphi_1)u^{(1)}(r_2, 0, g, \varphi_0, \varphi_1) + \\ &+ cu(r_2, 0, g, \varphi_0, \varphi_1) + u^{(1)}(r_1, 0, g, \varphi_0, \varphi_1)] + \\ &E \left[ c^2 u^{(1)}(r_1, 0, g, \varphi_0, \varphi_1) u^{(1)}(r_2, 0, g, \varphi_0, \varphi_1) \right]. \end{aligned}$$

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Let us consider the Dirichlet problem for the Helmholtz equation in a domain  $0 \leq x, y, z \leq 1$  :

$$\Delta u + cu = 0,$$
$$u|_{\Gamma} = \cos\left(x\sqrt{c/3}\right) \cos\left(y\sqrt{c/3}\right) \cos\left(z\sqrt{c/3}\right).$$

The exact solution is

$$u(x, y, z; c) = \cos\left(x\sqrt{c/3}\right) \cos\left(y\sqrt{c/3}\right) \cos\left(z\sqrt{c/3}\right).$$

For this problem  $c^* = 3\pi^2 \approx 29.609$ .

This problem solved in two ways. First the solution computed by standard random walks on spheres estimate

$$\eta_\varepsilon = u(r_N, c) \prod_{j=0}^{N-1} \frac{\sqrt{c} d_j}{\sin(\sqrt{c} d_j)}$$

at the point  $r = (x, y, z)$  where  $x = y = z = 0,9$ . The numerical results are presented in Table 1. Second the solution computed by probabilistic representation (??). Using Euler scheme with a constant time step  $\Delta t$ , approximations  $r_i$  of the trajectory for the random process  $\xi_t$  at the points of time  $i\Delta t$  was modeled statically. The solution computed by the following estimates

$$\zeta_1 = e^{cT} \varphi(r_N), \quad \zeta_2 = \left[ \sum_{i=0}^M \frac{c_0^i T^i}{i!} \right] e^{(c-c_0)T} \varphi(r_N).$$

Here  $T$  is the approximate moment of the first exit of the process  $r_j$  from on a boundary of the cub,  $r_N$  is approximate coordinate of an exit. The corresponding numerical results are presented in Table 2.



Table 1.




$c$	$N \times 10^{-6}$	$\varepsilon$	$-u(r)$	$-\tilde{u}(r)$	$ u(r) - \tilde{u}(r)  \pm \sqrt{\frac{Dn}{N}}$
50	16,7	$10^{-4}$	0.639	280.4	$281 \pm 5674$
50	1	$10^{-4}$	0.639	131162	$131162 \pm 131464$
35	16.7	$10^{-4}$	0.993	1.493	$0.5 \pm 0.260$
35	16.7	$10^{-2}$	0.993	0.337	$0.656 \pm 0.918$
30	9.8	$10^{-4}$	0.875	0.899	$0.023 \pm 0.018$
20	1	$10^{-4}$	0.3197	0.3252	$0.0055 \pm 0.004$
15	0.4	$10^{-4}$	0.078	0.0785	$0.0004 \pm 0.0003$

Table 2.




$c_0$	$c$	$N \times 10^6$	$\Delta t * 10^{-2}$	$-u(r)$	$-\tilde{u}(r)$	$ u(r) - \tilde{u}(r)  \pm$
–	30	10	1	0.875	0.790	$0.085 \pm 0.227$
–	35	1	1	0.993	5.385	$4.392 \pm 3.452$
5	30	1	1	0.875	1.143	$0.268 \pm 0.185$
5	30	$10^2$	1	0.875	1.086	$0.211 \pm 0.179$
10	35	$10^2$	$10^{-2}$	0.993	2.923	$1.93 \pm 0.516$

If the solution calculated by estimate  $\zeta_1$  then we put '–' in the first column of Table 2. We assume that  $M = 50$  when solution calculated by estimate  $\zeta_2$ .

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